# Beyond Closure Models: Learning Chaotic-Systems via Physics-Informed Neural Operators

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Small (numerical) error explodes along time!

• Long-term behavior/ statistics is of great practical importance in applications.



Climate Modeling



Aircraft Design

### **Chaotic Systems**

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Method 2: Coarse-grid simulations (CGS) + further modification ('error-correction')

### **Chaotic Systems**

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- Method 2: Coarse-grid simulations (CGS)+further modification ('error-correction')
  'Closure model'

## **Problem Setting**

#### [Notations]

- A nonlinear (and chaotic) dynamics  $\partial_t u = Lu, u \in H$
- $S_t$ : the semigroup (of the true dynamics)
- 'Trajectory':  $\{S_t u\}_{t \in R_{\geq 0}}$
- Attractor:  $\Omega \subset H \ s.t. \lim_{t \to \infty} dist(S(t)u, \Omega) = 0$ ,  $\forall u \in H$ .
- Invariant measure  $\mu^* \coloneqq \lim_{T \to \infty} \frac{1}{T} \int_0^T \delta_{S(t)u} dt$ , (independent of the initial u)
  - Measure in function space
  - Supported on  $\Omega$
- Statistics: For a functional *O*, the stat  $\langle O \rangle \coloneqq E_{u \sim \mu^*} O(u) = \int O(u) \mu^*(du)$

The Goal of this Project.

#### Outline

- (1) Review: data-driven closure models (& their inherent shortcomings)
- (2) New insight through Measure flow in function space
- (3) Our approach: Physics-informed Neural Operator

#### **Coarse-grid Simulations**

• Filtering operator  $F: u \to \overline{u}$ , e.g. spatial down sampling, Fourier-mode truncation. (Simulations with coarse grids)

Nonlinear dynamics:  $\partial_t u = Lu, u \in H$  (*H*: function space of interest) > Dynamics for  $\bar{u}$ :  $\partial_t \bar{u} = L\bar{u} + (FL - LF)u$ .

(nonlinear systemF and L does not commute)Simulating on low-res grids (the space of F(H))no access to  $\boldsymbol{u}$ .[Closure Model]: Evolve  $\partial_t \bar{u} = L\bar{u} + clos(\bar{u}; \theta)$ 

#### **Existing Data-driven Closure Models**

 $\succ$ [CGS]: Evolve  $\partial_t u = Lu + clos(u; \theta)$ 

Learned Model:

(1)Fix an ansatz  $clos(u; \theta)$ (2) Supervised loss:  $||(FL - LF)u - clos(Fu; \theta)||$ 

(3\*) More complex version:

 $clos: H(R^3) \to H(R^3)$ 

Training data  $\boldsymbol{u}$ : comes from costly Fully-Resolved simulations

#### **Existing Data-driven Closure Models**

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(1)Fix an ansatz clos(u; θ)
(2) Supervised loss: ||(FL – LF)u – clos(Fu; θ)||
(3\*) More complex version:
History-aware models:

 $clos: H(R^3) \rightarrow H(R^3)$ 

Training data  $\boldsymbol{u}$ : comes from costly Fully-Resolved simulations

$$clos: H(R^3 \times [0, t')) \to H(R^3)$$
  
$$u(x, t), x \in D, t \in [t_0 - t', t_0) \to clos(u)(x, t_0)$$

Stochastic models Different loss functions  $\succ [\text{Closure Model}]: \text{Evolve } \partial_t \bar{u} = L\bar{u} + clos(\bar{u}; \theta))$ 

- [Result 1] The target mapping clos(u) is not well-defined for all types of ansatz in the literature. (There are multiple possible outputs for the same input.)
- [Result 2] Following this scheme (directly adding a closure term), the model has to be trained with a large amount of fully-resolved data that suffices to compute the statistics.

We no longer need such a model!



(1) This is not a well-defined mapping.

(2) The resulting coarse-grid simulation might deviate from the filtered attractor.

(And the performance highly depends on the training data)

! We can only assign one moving direction in the reduced space F(H).

## **Theoretical Intuition**

 $\succ [FRS(ground truth)] \partial_t u = Lu$ 

 $\succ$ [CGS]: Evolve  $∂_t u = Lu + clos(u; θ)$ 

- Functions(states) *u* : infinite-dimensional particle systems in *H*.
- Only need to care about distributions (the invariant measure) Liouville flow/ Fokker-Planck equations for measure transformations.

 $\partial_t \rho = -\nabla \cdot (\boldsymbol{f} \rho)$ 

• The invariant measure is the solution to stationary Liouville/Fokker-Planck equation!

### Outline

- (1) Review: data-driven closure models (& their shortcomings)
- (2) Insight through Measure flow in function space
- (3) Our approach: Physics-informed Neural Operator

 $\succ$ [Previous Scheme]: Evolve  $\partial_t \bar{u} = L\bar{u} + clos(\bar{u};\theta)$ 

### Why do we have to adopt this form?

- New scheme: Evolve  $\partial_t \bar{u} = L(\theta) \bar{u}$
- Intuition: the error-correcting term is implicitly involved in the simulation.

#### (Fourier) Neural Operator



(a) The full architecture of neural operator: start from input a. 1. Lift to a higher dimension channel space by a neural network P. 2. Apply four layers of integral operators and activation functions. 3. Project back to the target dimension by a neural network Q. Output u. (b) Fourier layers: Start from input v. On top: apply the Fourier transform  $\mathcal{F}$ ; a linear transform R on the lower Fourier modes and filters out the higher modes; then apply the inverse Fourier transform  $\mathcal{F}^{-1}$ . On the bottom: apply a local linear transform W.

Figure 2: top: The architecture of the neural operators; bottom: Fourier layer.

 $\mathcal{G}_{FNO} := \mathcal{Q} \circ (W_L + \mathcal{K}_L) \circ \cdots \circ \sigma (W_1 + \mathcal{K}_1) \circ \mathcal{P},$ 

The input can be either course-grid or fine-grid.

- New scheme: Evolve  $\partial_t \bar{u} = L(\theta) \bar{u}$
- Intuition: the error-correcting term is implicitly involved in the simulation.
   [Method]
- Physics-informed Neural Operator

Much fewer DNS data needed

Ansatz for  $L(\theta)$ Merit: resolution invarianc\_



$$J_{pde}(\theta; \mathfrak{D}) = \frac{1}{|\mathfrak{D}|} \sum_{i \in \mathfrak{D}} \|(\partial_t - \mathcal{A}) \mathcal{G}_{\theta} u_{0i}(x)\|_{L^2(\Omega \times [0,h])}$$

• Provable convergence guarantee for statistics.

**Theorem 3.1.** For any h > 0, denote  $\hat{\mu}_{h,\theta} := \lim_{N \to \infty} \frac{1}{N} \sum_{n=1}^{N} \delta_{\mathcal{G}_{\theta}^{n}v_{0}(x)}$ , any  $v_{0}(x)$  with  $x \in D'$ . For any  $\epsilon > 0$ , there exists  $\delta > 0$  s.t. as long as  $\|(\mathcal{G}_{\theta}u)(\cdot, h) - S(h)u\|_{\mathcal{H}} < \delta, \forall u \in \mathcal{H}$ , we have  $\mathcal{W}_{\mathcal{H}}(\hat{\mu}_{h,\theta}, \mathcal{F}_{\#}\mu^{*}) < \epsilon$ , where  $\mathcal{W}_{\mathcal{H}}$  is a generalization of Wasserstein distance in function space.

### **Summary**

Method	Optimal statistics	High-res. training data DNS Snapshots   Trajs.		Complexity
Fully-resolved Simulation, e.g., DNS [33, 34]	✓ ×	-	$Re^{3.52}$ $Re^{2.48}$	
Coarse-grid Simulation, e.g., LES [55, 54]	^	-	P. 14	ne
Single-state model [28]	×	24000	8	$Re^{2.48}$
History-aware model[35]	×	250000	50	$Re^{2.48}$
Latent Neural SDE[32]	×	179200	28	$\frac{1}{\delta t} Re^{1.86}$
Physics-Informed Operator Learning (Ours)	1	110	1	$Re^{1.86}$

Re: Reynolds number

- 1. Require large number of DNS data (which are not available usually).
- 2. Require a coarse-grid solver (can be even faster).
- 3. Cannot give the optimal estimations of statistics in ideal case, i.e. perfect training.

#### **1D Kuramoto–Sivashinsky (toy test example)**

 $\partial_t u + u \partial_x u + \partial_{xx} u + \nu \partial_{xxxx} u = 0, \quad (x, t) \in [0, 6\pi] \times \mathbb{R}_+,$ 

- Viscosity: 0.01 L:  $6\pi$
- DNS: 1024 spatial grid
- LES: 128 spatial grid
- Note: baseline methods are trained with the same number of DNS data as ours.

Method	Avg. Eng.	Max Eng.	Avg. Cor.	Max Cor.	Velocity	Avg. TV	Max TV
CGS (No closure)	12.5169%	77.8223%	13.1275%	80.5793%	0.0282	0.0398	0.2097
Eddy-Viscosity [57]	7.6400%	48.3684%	8.7583%	56.5878%	0.0276	0.0282	0.1462
Single-state [28]	12.5323%	78.6410%	13.1052%	81.2461%	0.0280	0.0410	0.2111
Our Method	7.4776%	20.4176%	7.8706%	22.7046%	0.0284	0.0272	0.0849

Energy Spectrum Correlation Velocity Total Variation of each mode

#### **2D Kolmogorov Flow (toy test example)**

$$\partial_t \mathbf{u} = -(\mathbf{u} \cdot \nabla)\mathbf{u} - \nabla p + \frac{1}{Re}\Delta \mathbf{u} + (\sin(4y), 0)^T, \quad \nabla \cdot \mathbf{u} = 0, \quad (x, y, t) \in [0, L]^2 \times \mathbb{R}_+,$$

- Reynolds number: 100 L:  $2\pi$
- DNS: 128\*128 spatial grid
- LES: 16\*16 spatial grid
- Note: baseline methods are trained with the same number of DNS data as ours.

Method	Avg. Eng.	Max Eng.	Vorticity	Avg. TV	Max TV	Variance
CGS (No closure) Smagorinsky [14] Single-state [28]	178.4651% 52.9511% 205.3709%	404.9923% 120.0723% 487.3957%	0.1512 0.0483 0.1648	0.4914 0.2423 0.5137	0.8367 0.9195 0.8490	253.4234% 20.1740% 298.2027%
Our Method	5.3276%	8.9188%	0.0091	0.0726	0.2572	2.8666%

Energy Spectrum Vorticity Total Variation of each mode



#### **Summary**

Takeaway: We need to ensure a well-defined target mapping before any 'learning'. See more details in https://arxiv.org/pdf/2408.05177

# **Thanks!**